# Generalization of the Titchmarsh's theorem for the second Hankel-Clifford transformation 

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#### Abstract

Using a generalized translation operator, we obtain an analogue of Titchmarsh's theorem for the second Hankel-Clifford transformation for functions satisfying the second Hankel-Clifford Lipschitz condition in the space $L_{\mu}^{2}\left((0,+\infty), x^{\mu}\right)$.


Keywords: Generalized translation operator, Second Hankel-Clifford transformation, Second Hankel-Clifford Lipschitz class
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## 1 Introduction and preliminaries

The theorem 85 in [15, Titchmarsh characterized the set of functions in $L^{2}(\mathbb{R})$ satisfying the Cauchy Lipschitz class by means of an asymptotic estimate growth of the norm of their Fourier transform, namely we have

Theorem 1.1. ([15], Theorem 85) Let $\alpha \in(0,1)$ and assume that $f \in L^{2}(\mathbb{R})$. Then the following are equivalents

1. $\|f(x+h)-f(x)\|_{L^{2}(\mathbb{R})}=O\left(h^{\alpha}\right)$ as $h \longrightarrow 0$
2. $\int_{|\lambda| \geq s}|\hat{f}(\lambda)|^{2} d \lambda=O\left(s^{-2 \alpha}\right)$,
where $\hat{f}$ stands for the Fourier transform of $f$.
In this paper we obtain an analogue of this theorem 1.1 for the Second Hankel-Clifford transformation. There are many analogues of this result: for the Fourier transform, for the Jacobi transform, for the Fourier transform on the group of p-Adic Numbers, For the Fourier-Walsh transform, for the generalized Dunkl transform, for the generalized Bessel transform etc (see, for exemple [3, 4, 5, 6, 12, 13]).

We briefly overview the theory of second Hankel-Clifford transformation and related harmonic analysis (see [10, 11, 14).

We define the space $L_{\mu}^{p}=L_{\mu}^{p}((0,+\infty)), \quad 1 \leq p<\infty$ and $\mu \geq 0$, as the space of all those real-valued measurable functions $f$ on $(0,+\infty)$, such that

[^0]$$
\|f\|_{L_{\mu}^{p}}=\left(\int_{0}^{\infty}|f(x)|^{p} x^{\mu} d x\right)^{1 / p}<\infty
$$

The Bessel-Clifford function of the first kind of order $\mu \geq 0$ (See [7]).

$$
c_{\mu}(x)=\sum_{k=0}^{+\infty} \frac{(-1)^{k} x^{k}}{k!\Gamma(\mu+k+1)},
$$

is a solution of the differential equation

$$
x y^{\prime \prime}+(\mu+1) y^{\prime}+y=0,
$$

and we have

$$
\begin{equation*}
c_{\mu}(x)=x^{-\frac{\mu}{2}} J_{\mu}(2 \sqrt{x}), \tag{1.1}
\end{equation*}
$$

where $J_{\mu}$ the Bessel function of first kind.

For $f \in L_{\mu}^{1}$. Hayek [9] introduced the second Hankel-Clifford transformation by

$$
h_{2, \mu}(f)(\lambda)=\int_{0}^{+\infty} c_{\mu}(\lambda x) f(x) x^{\mu} d x
$$

and its inversion formula is defined by

$$
f(x)=\int_{0}^{+\infty} c_{\mu}(\lambda x) h_{2, \mu}(f)(\lambda) \lambda^{\mu} d \lambda
$$

The corresponding Parseval's equality now takes the form [11]

$$
\int_{0}^{+\infty} f(x) g(x) x^{\mu} d x=\int_{0}^{+\infty} F_{2}(\lambda) G_{2}(\lambda) \lambda^{\mu} d \lambda
$$

where $F_{2}(\lambda)=h_{2, \mu}(f)(\lambda)$ and $G_{2}(\lambda)=h_{2, \mu}(g)(\lambda)$.
i.e,. For $f \in L_{\mu}^{2}$, we have

$$
\|f\|_{L_{\mu}^{2}}=\left\|h_{2, \mu}(f)\right\|_{L_{\mu}^{2}} .
$$

Let $\Delta=\Delta(x, y, z)$ be the area of triangle with sides $x, y, z($ see [8, 16]). Set

$$
D_{\mu}(x, y, z)=\frac{\Delta^{2 \mu+1}}{2^{2 \mu}(x y z)^{\mu} \Gamma\left(\mu+\frac{1}{2}\right) \sqrt{\pi}}
$$

If $\Delta$ exists and zero otherwise. We note that $D_{\mu}(x, y, z) \geq 0$ and it is symmetric in $x, y, z$.
The generalized translation operator value of $f \in L_{\mu}^{2}$ is defined by

$$
T_{h}(f)(x)=\int_{0}^{+\infty} f(z) D_{\mu}(h, x, z) z^{\mu} d z, 0<x, h<\infty
$$

From lemma 1.3 in [14], we have

$$
\begin{equation*}
h_{2, \mu}\left(T_{h}(f)\right)(\lambda)=c_{\mu}(\lambda h) h_{2, \mu}(f)(\lambda) \tag{1.2}
\end{equation*}
$$

where $f \in L_{\mu}^{2}$.
For $\mu \geq-\frac{1}{2}$, we introduce the normalized spherical Bessel function $j_{\mu}$ defined by

$$
\begin{equation*}
j_{\mu}(x)=\frac{2^{\mu} \Gamma(\mu+1) J_{\mu}(x)}{x^{\mu}} \tag{1.3}
\end{equation*}
$$

From [1] we have the following lemma:

Lemma 1.2. Let $\mu \geq-\frac{1}{2}$. The following inequalities hold

1. $\left|j_{\mu}(x)\right| \leq 1$
2. $1-j_{\mu}(x)=O\left(x^{2}\right) ; \quad 0 \leq x \leq 1$
3. $\sqrt{x} J_{\mu}(x)=O(1)$.

Lemma 1.3. The following inequality is true

$$
\left|1-j_{\mu}(x)\right| \geq c
$$

with $|x| \geq 1$, where $c>0$ is certain constant.

Proof . Analog of lemma 2.9 in [2].
It follows from (1.1) and (1.3) that

$$
c_{\mu}(x)=\frac{1}{\Gamma(\mu+1)} j_{\mu}(2 \sqrt{x}) .
$$

## 2 Main result

In this section we give the main result of this paper. We need first to define the second Hankel-Clifford Lipschitz class.

Definition 2.1. Let $\alpha \in(0,1)$. A function $f \in L_{\mu}^{2}$ is said to be in the second Hankel-Clifford Lipschitz class, denoted by $\operatorname{Lip}(\alpha, 2, \mu)$, If

$$
\left\|T_{h} f(x)-\frac{1}{\Gamma(\mu+1)} f(x)\right\|_{L_{\mu}^{2}}=O\left(h^{\alpha}\right) \quad \text { as } h \longrightarrow 0 .
$$

Our main result is the next theorem

Theorem 2.2. Let $f \in L_{\mu}^{2}$. Then the following are equivalent:

1. $f \in \operatorname{Lip}(\alpha, 2, \mu)$
2. $\int_{N}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(N^{-2 \alpha}\right) \quad$ as $N \longrightarrow+\infty$

Proof . 1) $\Longrightarrow 2)$ Let $f \in L_{\mu}^{2}$. It follows from (1.2) and (1.4) that

$$
\begin{aligned}
h_{2, \mu}\left(T_{h} f-\frac{1}{\Gamma(\mu+1)} f\right)(\lambda) & =\left(C_{\mu}(\lambda h)-\frac{1}{\Gamma(\mu+1)}\right) h_{2, \mu}(f)(\lambda) \\
& =\frac{1}{\Gamma(\mu+1)}\left(j_{\mu}(2 \sqrt{\lambda h})-1\right) h_{2, \mu}(f)(\lambda)
\end{aligned}
$$

then, using the Parseval's identity, we have

$$
\left\|T_{h} f-\frac{1}{\Gamma(\mu+1)} f\right\|_{L_{\mu}^{2}}^{2}=\frac{1}{\Gamma(\mu+1)} \int_{0}^{+\infty}\left|1-j_{\mu}(2 \sqrt{\lambda h})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda .
$$

Assume that $f \in \operatorname{Lip}(\alpha, 2, \mu)$. Then we have

$$
\left\|T_{h} f-\frac{1}{\Gamma(\mu+1)} f\right\|_{L_{\mu}^{2}}=O\left(h^{\alpha}\right) \quad \text { as } h \longrightarrow 0
$$

If $\lambda \in\left[\frac{1}{4 h}, \frac{2}{4 h}\right]$, then $2 \sqrt{\lambda h} \geq 1$. From lemme 1.3 we obtain

$$
1 \leq \frac{1}{c^{2}}\left|1-j_{\mu}(2 \sqrt{\lambda h})\right|^{2}
$$

i.e.,

$$
\frac{1}{\Gamma(\mu+1)} \leq \frac{1}{c^{2} \Gamma(\mu+1)}\left|1-j_{\mu}(2 \sqrt{\lambda h})\right|^{2}
$$

Then

$$
\begin{aligned}
\frac{1}{\Gamma(\mu+1)} \int_{\frac{1}{4 h}}^{\frac{2}{4 h}}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda & \leq \frac{1}{c^{2} \Gamma(\mu+1)} \int_{\frac{1}{4 h}}^{\frac{2}{4 h}}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \\
& \leq \frac{1}{c^{2} \Gamma(\mu+1)} \int_{0}^{+\infty}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \\
& =O\left(h^{2 \alpha}\right)
\end{aligned}
$$

we conclude that

$$
\int_{N}^{2 N}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(N^{-2 \alpha}\right) \quad \text { as } N \longrightarrow+\infty
$$

Thus there exists $C_{1}>0$ such that

$$
\int_{N}^{2 N}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \leq C_{1} N^{-2 \alpha}
$$

So that

$$
\begin{aligned}
\int_{N}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda & =\left(\int_{N}^{2 N}+\int_{2 N}^{4 N}+\int_{4 N}^{8 N}+\ldots .\right)\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \\
& \leq C_{1}\left(N^{-2 \alpha}+(2 N)^{-2 \alpha}+(4 N)^{-2 \alpha}+\ldots\right) \\
& \leq C_{1} N^{-2 \alpha}\left(1+2^{-2 \alpha}+\left(2^{-2 \alpha}\right)^{2}+\left(2^{-2 \alpha}\right)^{3}+\ldots\right) \\
& \leq C_{1} K_{\alpha} N^{-2 \alpha},
\end{aligned}
$$

where $K_{\alpha}=\left(1-2^{-2 \alpha}\right)^{-1}$ since $2^{-2 \alpha}<1$.
This proves that

$$
\int_{N}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(N^{-2 \alpha}\right) \text { as } N \longrightarrow+\infty .
$$

$2) \Longrightarrow 1)$ Suppose now that

$$
\int_{N}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(N^{-2 \alpha}\right) \text { as } N \longrightarrow+\infty
$$

we have to show that

$$
\frac{1}{\Gamma(\mu+1)} \int_{0}^{+\infty}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(h^{2 \alpha}\right) \text { as } h \longrightarrow 0
$$

We write

$$
\int_{0}^{+\infty}\left|1-j_{\mu}(2 \sqrt{\lambda h})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=I_{1}+I_{2}
$$

where

$$
I_{1}=\int_{0}^{\frac{1}{4 h}}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda,
$$

and

$$
I_{2}=\int_{\frac{1}{4 h}}^{+\infty}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda
$$

From (1) of lemma 1.2, we have

$$
\begin{aligned}
I_{2} & =\int_{\frac{1}{4 h}}^{+\infty}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \\
& \leq 4 \int_{\frac{1}{4 h}}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda \\
& =O\left(h^{2 \alpha}\right) \text { as } h \longrightarrow 0 .
\end{aligned}
$$

Then

$$
\frac{1}{\Gamma(\mu+1)} \int_{\frac{1}{4 h}}^{+\infty}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(h^{2 \alpha}\right)
$$

Set

$$
\psi(x)=\int_{x}^{+\infty}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda
$$

We know from (2) of lemma 1.2 that $1-j_{\mu}(2 \sqrt{h \lambda})=O(\lambda h)$ for $0 \leq 2 \sqrt{\lambda h} \leq 1$. Thus there exists $C_{2}>0$ such that $\left|1-j_{\mu}(2 \sqrt{h \lambda})\right| \leq C_{2} \lambda h$ for $0 \leq 2 \sqrt{\lambda h} \leq 1$. Then

$$
I_{1} \leq-C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x^{2} \psi^{\prime}(x) d x
$$

An integration by parts yields

$$
\begin{aligned}
I_{1} & \leq-C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x^{2} \psi^{\prime}(x) d x \\
& \leq-C_{2} \psi\left(\frac{1}{4 h}\right)+2 C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x^{2} \psi(x) d x \\
& \leq 2 C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x \psi(x) d x \\
& \leq 2 C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x x^{-2 \alpha} d x \\
& \left.\leq 2 C_{2} h^{2} \int_{0}^{\frac{1}{4 h}} x^{1-2 \alpha} d x \quad \text { (the integral exists since } \alpha<1\right) \\
& \leq C_{2} K h^{2 \alpha} .
\end{aligned}
$$

where $K$ is a positive constant. Then

$$
\frac{1}{\Gamma(\mu+1)} \int_{0}^{\frac{1}{4 h}}\left|1-j_{\mu}(2 \sqrt{h \lambda})\right|^{2}\left|h_{2, \mu}(f)(\lambda)\right|^{2} \lambda^{\mu} d \lambda=O\left(h^{2 \alpha}\right)
$$

and this ends the proof.

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