# Multiplicity analysis of positive weak solutions in a quasi-linear Dirichlet problem inspired by Kirchhoff-type phenomena 

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#### Abstract

The main focus of this paper lies in investigating the existence of infinitely many positive weak solutions for the following elliptic-Kirchhoff equation with Dirichlet boundary condition $$
\begin{cases}-\sum_{\substack{i=1}}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right)=f(x, u) & \text { in } \Omega, \\ u=0 & \text { on } \partial \Omega .\end{cases}
$$


The methodology adopted revolves around the technical approach utilizing the direct variational method within the framework of anisotropic variable exponent Sobolev spaces.

Keywords: Nonlinear elliptic equations, Variational methods applied on PDEs, Positive solutions to PDEs
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## 1 Introduction

Over the past years, differential equations have been a focal point of research, owing to their extensive practical implications and widespread use in numerous fields.

Proposed by Kirchhoff [20], the Kirchhoff differential equations offer an extension to D'Alembert's wave equation, accommodating the effects of string length changes during vibrations

$$
\begin{equation*}
\rho \frac{\partial^{2} u}{\partial t^{2}}-\left(\frac{P_{0}}{h}+\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2} d x\right) \frac{\partial^{2} u}{\partial x^{2}}=0 \tag{1.1}
\end{equation*}
$$

in this context, $L$ denotes the length of the chord, $h$ represents the area of the cross-section, $E$ stands for the Young's modulus of the material, denotes the density, and $P_{0}$ corresponds to the initial tension. The Kirchhoff equation 1.1

[^0]exhibits a unique characteristic in its inclusion of a non-local coefficient $\frac{P_{0}}{h}+\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2} d x$, which is dependent on the average $\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2}$ of the kinetic energy $\frac{1}{2}\left|\frac{\partial u}{\partial x}\right|^{2}$ within the interval $[0, L]$. Thus, the equation loses its property of being a point-wise identity. See also [6, 19, 32] for related topics.

In recent times, numerous mathematicians, physicists, and engineers have shown a keen interest in anisotropic variable exponent Sobolev spaces. The motivation behind this stems from the crucial role these spaces play in modeling real-world phenomena, including electrorheological and image restoration, magnetorheological fluids, and elastic materials, (look at, for example [5, 8, 10, 30, 33, 34, 35, 36]).

In the present paper we study the existence of positive solutions of the nonhomogeneous anisotropic Kirchhoff problem

$$
\begin{cases}-\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right)=f(x, u) & \text { in } \Omega  \tag{1.2}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

where $\Omega \subset \mathbb{R}^{N},(N>3)$ represents a bounded domain with a smooth boundary $\partial \Omega$, and $p_{i}, i=1, \ldots N$ are continuous functions. Additionally, for each $i=1, \ldots N, M_{i}$ and $f$ are continuous functions which satisfies some conditions detailed in Section 3.

The differential operator

$$
\begin{equation*}
\Delta_{\vec{p}(\cdot)} u=\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right), \tag{1.3}
\end{equation*}
$$

involved in problem (1.2) is an anisotropic variable exponent $\vec{p}(\cdot)$-Laplace operator which represents an extension of the operator

$$
\begin{equation*}
\Delta_{p(\cdot)} u=\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u\right) \tag{1.4}
\end{equation*}
$$

The $p(\cdot)$-Laplacian operator, obtained by setting each $p_{i}(x)$ to be equal to $p(x)$ for $i=1, \ldots N$, serves as a natural extension of the isotropic $p$-Laplacian operator

$$
\begin{equation*}
\Delta_{p} u=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right) \tag{1.5}
\end{equation*}
$$

where $p>1$ denotes a real constant. In the classical Sobolev spaces, F. J. S. A. Corrêa, R. G. Nascimento [12] have established the existence of solutions for problem 1.2 in this particular case $p$-Kirchhoff-type equation, for additional results, refer to [22, 28, 29, 31.

In the Sobolev variable exponent setting, G. Dai and D. Liu [13] has analyzed the problem 1.2 in the context of $p(x)$-Kirchhoff-type equation, see also [2, 9, 11, 14, 18] for related topics.

The investigation of problem (1.2) in anisotropic variable exponent Sobolev spaces has been previously addressed by other researchers (see [7, 16, 25]). However, our study stands apart due to the entirely distinct hypotheses adopted, which subsequently lead to different and novel findings.

The shift from a variable exponent to an anisotropic variable exponent inevitably introduces fresh complexities. To tackle these challenges, we adopt a combined approach, utilizing traditional methodologies alongside modern techniques specifically designed for handling problems of anisotropic nature with variable exponents. The organization of this paper is as follows: In Section 2, we provide an introduction to anisotropic variable exponent Sobolev spaces, laying the necessary groundwork for the subsequent analysis. Section 3 is dedicated to presenting the assumptions under which our problem yields positive solutions, accompanied by an illustrative example.

## 2 Preliminary

Let $\Omega$ denote a smooth bounded domain in $\mathbb{R}^{N}$, where we introduce the following definitions:

$$
\mathcal{C}_{+}(\bar{\Omega})=\left\{p \in \mathcal{C}(\bar{\Omega}) \quad \text { such that } \quad 1<p^{-} \leq p^{+}<\infty\right\}
$$

where

$$
p^{-}=\operatorname{ess} \inf \{p(x): x \in \bar{\Omega}\} \quad \text { and } \quad p^{+}=\operatorname{ess} \sup \{p(x): x \in \bar{\Omega}\}
$$

For any $p \in \mathcal{C}_{+}(\bar{\Omega})$, we introduce the Lebesgue space with variable exponent $L^{p(\cdot)}(\Omega)$, which encompasses all measurable functions $u: \Omega \longmapsto \mathbb{R}$ such that the convex modular

$$
\rho_{p(\cdot)}(u):=\int_{\Omega}|u|^{p(x)} d x
$$

remains finite. Consequently, we define the norm

$$
\|u\|_{L^{p(\cdot)}(\Omega)}=\|u\|_{p(\cdot)}=\inf \left\{\lambda>0: \rho_{p(\cdot)}\left(\frac{u}{\lambda}\right) \leq 1\right\}
$$

as the Luxemburg norm in $L^{p(\cdot)}(\Omega)$. As a separable Banach space, $\left(L^{p(\cdot)}(\Omega),|\cdot|_{p(\cdot)}\right)$ exhibits desirable properties. Additionally, the space $L^{p(\cdot)}(\Omega)$ is uniformly convex, making it reflexive, and its dual space is isomorphic to $L^{p^{\prime}(\cdot)}(\Omega)$, with $\frac{1}{p(\cdot)}+\frac{1}{p^{\prime}(\cdot)}=1$. Lastly, we arrive at the following Hölder-type inequality.

$$
\begin{equation*}
\left|\int_{\Omega} u(x) v(x) d x\right| \leq\left(\frac{1}{p^{-}}+\frac{1}{\left(p^{-}\right)^{\prime}}\right)\|u\|_{p(\cdot)}\|v\|_{p^{\prime}(\cdot)} \quad \text { for all } u \in L^{p(\cdot)}(\Omega) \text { and } v \in L^{p^{\prime}(\cdot)}(\Omega) \tag{2.1}
\end{equation*}
$$

The modular $\rho_{p(\cdot)}$ of the space $L^{p(\cdot)}(\Omega)$ assumes a crucial role in handling the generalized Lebesgue spaces. The ensuing result is presented:

Proposition 2.1. (See [15]). Considering $u_{n}, u \in L^{p(\cdot)}(\Omega)$, with $p^{+}<+\infty$, we observe the subsequent properties:

1. If $\|u\|_{p(\cdot)}>1$, then $\|u\|_{p(\cdot)}^{p^{-}}<\rho_{p(\cdot)}(u)<\|u\|_{p(\cdot)}^{p^{+}}$.
2. For $\|u\|_{p(\cdot)}<1$, we have $\|u\|_{p(\cdot)}^{p^{+}}<\rho_{p(\cdot)}(u)<\|u\|_{p(\cdot)}^{p^{-}}$.
3. The condition $\|u\|_{p(\cdot)}<1$ (respectively $=1 ;>1$ ) is equivalent to $\rho_{p(\cdot)}(u)<1$ (respectively $=1 ;>1$ ).
4. When $\left\|u_{n}\right\|_{p(\cdot)} \rightarrow 0$ (respectively $\rightarrow+\infty$ ), it implies $\rho_{p(\cdot)}\left(u_{n}\right) \rightarrow 0$ (respectively $\rightarrow+\infty$ ).
5. Lastly, we have $\rho_{p(\cdot)}\left(\frac{u}{\|u\|_{p(\cdot)}}\right)=1$.

The definition of $W_{0}^{1, p(\cdot)}(\Omega)$ involves taking the closure of $\mathcal{C}_{0}^{\infty}(\Omega)$ in $W^{1, p(\cdot)}(\Omega)$, and

$$
p^{*}(x)=\left\{\begin{array}{cc}
\frac{N p(x)}{N-p(x)} & \text { for } p(x)<N \\
\infty & \text { for } p(x) \geq N
\end{array}\right.
$$

Proposition 2.2. (See [15).
(i) For $1<p^{-} \leq p^{+}<\infty$, both $W^{1, p(\cdot)}(\Omega)$ and $W_{0}^{1, p(\cdot)}(\Omega)$ are separable and reflexive Banach spaces.
(ii) If $q(x) \in C_{+}(\bar{\Omega})$ and $q(x)<p^{*}(x)$ holds true for each $x \in \Omega$, then the embedding $W_{0}^{1, p(\cdot)}(\Omega) \hookrightarrow \hookrightarrow L^{q(\cdot)}(\Omega)$ is continuous and compact.

We now introduce the anisotropic Sobolev space with variable exponent, which serves as the foundation for studying our main problem. Consider $N$ variable exponents $p_{1}(\cdot), \ldots, p_{N}(\cdot)$ belonging to $\mathcal{C}_{+}(\bar{\Omega})$. We use the notation

$$
\vec{p}(\cdot)=\left\{p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\} \text { and } D^{i} u=\frac{\partial u}{\partial x_{i}} \quad \text { for } i=1, \ldots, N
$$

and we set it for all $x \in \bar{\Omega}$,

$$
p_{M}(\cdot)=\max \left\{p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\} \text { and } p_{m}(\cdot)=\min \left\{p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\} .
$$

The following notations are introduced:

$$
\begin{equation*}
\underline{p}=\min \left\{p_{1}^{-}, p_{2}^{-}, \ldots, p_{N}^{-}\right\}, \quad \underline{p}^{+}=\max \left\{p_{1}^{-}, p_{2}^{-}, \ldots, p_{N}^{-}\right\}, \quad \bar{p}=\max \left\{p_{1}^{+}, p_{2}^{+}, \ldots, p_{N}^{+}\right\}, \tag{2.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\underline{p}^{*}=\frac{N}{\sum_{i=1}^{N} \frac{1}{p_{i}^{-}}-1}, \quad \underline{p}, \infty=\max \left\{\underline{p}^{*}, \underline{p}^{+}\right\} . \tag{2.3}
\end{equation*}
$$

In the context of this paper, we make the assumption that

$$
\begin{equation*}
\sum_{i=1}^{N} \frac{1}{p_{i}^{-}}>1 \tag{2.4}
\end{equation*}
$$

The definition of the anisotropic variable exponent Sobolev space $W^{1, \vec{p}(\cdot)}(\Omega)$ is as follows:

$$
W^{1, \vec{p}(\cdot)}(\Omega)=\left\{D^{i} u \in L^{p_{i}(\cdot)}(\Omega), \quad i=1,2, \ldots, N\right\}
$$

equipped with the norm

$$
\begin{equation*}
\|u\|_{W^{1, \vec{p}(\cdot)}(\Omega)}=\|u\|_{1, \vec{p}(\cdot)}=\sum_{i=1}^{N}\left\|D^{i} u\right\|_{L^{p_{i}(\cdot)}(\Omega)}, \tag{2.5}
\end{equation*}
$$

Furthermore, $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is defined as the closure of $\mathcal{C}_{0}^{\infty}(\Omega)$ in $W^{1, \vec{p}(\cdot)}(\Omega)$ under the norm 2.5. The dual space of $W 0^{1, \vec{p}(\cdot)}(\Omega)$ is denoted as $W^{-1, \vec{p}(\cdot)^{\prime}}(\Omega)$, where $\overrightarrow{p^{\prime}}(x)=p_{0}^{\prime}(x), \ldots, p_{N}^{\prime}(x)$, satisfying $\frac{1}{p_{i}^{\prime}(x)}+\frac{1}{p_{i}(x)}=1$ (see [26, 27]) for the constant exponent case). The reflexivity of the Banach space $\left(W_{0}^{1, \vec{p}(\cdot)}(\Omega),|u|_{1, \vec{p}(\cdot)}\right)$ has been established in [24]. For a more comprehensive treatment of anisotropic variable exponent Sobolev spaces, researchers may delve into [1, 3, 17, 21, 24].

Proposition 2.3. (See [4, 23]). The bounded domain $\Omega \subset \mathbb{R}^{N}$, with a smooth boundary and $N>3$, satisfies relation (2.4).

1. Considering any $q \in \mathcal{C}+(\bar{\Omega})$ satisfying the condition $1<q(x)<\underline{p}, \infty$, for all $x \in \bar{\Omega}$, then

$$
W_{0}^{1, \vec{p}(\cdot)}(\Omega) \hookrightarrow \hookrightarrow L^{q(\cdot)}(\Omega)
$$

2. Assume that $\underline{p}>N$ then

$$
W_{0}^{1, \vec{p}(\cdot)}(\Omega) \hookrightarrow \hookrightarrow C^{0}(\bar{\Omega})
$$

## 3 Fundamental assumptions and main results

For the entirety of this paper, we make the assumption that the following set of conditions is satisfied:
Assume that $f: \Omega \times \mathbb{R} \longmapsto \mathbb{R}$ is Carathéodory functions satisfying the following condition
(H1) There exists a constant $\tau>0$ such that $\sup _{t \in[0, \tau]} f(\cdot, t) \in L^{\infty}(\Omega)$.
(H2) Suppose that $\left(a_{n}\right)_{n}$ and $\left(b_{n}\right)_{n}$ be two positive sequences such that

$$
0<a_{n}<b_{n}, \lim _{n \rightarrow \infty} b_{n}=0, \text { and } \int_{0}^{a_{n}} f(x, s) d s=\sup _{t \in\left[a_{n}, b_{n}\right]} \int_{0}^{t} f(x, s) d s \text { for almost all } x \in \Omega \text { and } n \in \mathbb{N} .
$$

(H3) There is a sequence $\left(\vartheta_{n}\right)_{n}$, which is a subset of the interval $\left[0, b_{n}\right]$, such that

$$
e s s \inf _{\Omega} \int_{0}^{\vartheta_{n}} f(x, s) d s>0
$$

For the function $M_{i}, i=1, \ldots N$, we set forth the subsequent assumptions.
(H4) $M_{i}$ is a differentiable on $\mathbb{R}^{+}$and there is positive constant $m$ such that

$$
M_{i}(t) \geq m \text { for all } t \geq 0
$$

Functionals are defined for any $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ as follows:

$$
\begin{equation*}
\Phi(u)=\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} F(x, u) d x \tag{3.1}
\end{equation*}
$$

where $F(x, t)=\int_{0}^{t} f(x, s) d s$ and $\widehat{M}_{i}(t)=\int_{0}^{t} M_{i}(s) d s$.
Definition 3.1. For any measurable function $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ to be considered a weak solution of the elliptic problem (1.2), it must satisfy the condition that, for all $v \in W_{0}^{1, \vec{p} \cdot)}(\Omega)$,

$$
\begin{equation*}
\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \int_{\Omega} \sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} d x=\int_{\Omega} f(x, u) v(x) d x \tag{3.2}
\end{equation*}
$$

It is easy to see that $\Phi \in \mathcal{C}^{1}\left(W_{0}^{1, \vec{p} \cdot \cdot}(\Omega), \mathbb{R}\right)$ (see [7, 25] ), and the function $u \in W_{0}^{1, \vec{p} \cdot \cdot)}(\Omega)$ is deemed a weak solution of (1.2) if and only if it corresponds to a critical point of the functional $\Phi$.

Considering our assumptions on $f$, we can find positive constants $k$ and $\tau$ such that $|f(\cdot, t)| \leq k$ for every $0 \leq \tau \leq t$ and almost every $x \in \Omega$. Without any loss of generality, we can suppose that $b_{n} \leq \tau$ for every $n \in \mathbb{N}$. Let's proceed by defining

$$
g(\cdot, t)= \begin{cases}0 & \text { if } t \leq 0  \tag{3.3}\\ f(\cdot, t) & \text { if } 0<t \leq \tau \\ f(\cdot, \tau) & \text { if } t>\tau\end{cases}
$$

Hence, we have

$$
\begin{equation*}
|g(\cdot, t)| \leq k \tag{3.4}
\end{equation*}
$$

for almost every $x \in \Omega$ and every $t \in \mathbb{R}$. Next, we take into account the following problem

$$
\begin{cases}-\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right)=g(x, u) & \text { in } \Omega  \tag{3.5}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

We can identify the weak solutions of (3.5) as the critical points of the functional

$$
\begin{equation*}
\Psi(u)=\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} G(x, u) d x \tag{3.6}
\end{equation*}
$$

where $G(x, t)=\int_{0}^{t} g(x, s) d s$. By (3.4), it is clear that $\Psi$ is well defined and Gâteaux differentiable in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ (see [7, [25]). For every fixed $n \in \mathbb{N}$, we define

$$
\begin{equation*}
K_{n}(u)=\left\{u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega): 0<u(x) \leq b_{n} \text { a.e. } \Omega\right\} \tag{3.7}
\end{equation*}
$$

Having established the necessary groundwork, we can now present the main findings of this paper.
Theorem 3.2. Assume assumptions (H1) (H4) hold true and $f(\cdot, 0)=0$. Then, there exists a sequence $\left(u_{n}\right)_{n} \subset$ $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ of positive, homoclinic weak solutions of 1.2$)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \Psi\left(u_{n}\right)=0 \text { and } \lim _{n \rightarrow+\infty}\left\|u_{n}\right\|_{1, \vec{p}(\cdot)}=0 \tag{3.8}
\end{equation*}
$$

Theorem 3.3. To enhance the organization and clarity, we divided the proof into three steps.

Step 1 :. Auxiliary lemmas.
Lemma 3.4. Assume assumptions (H1), (3.4) and (H4) are satisfied. Then, the functionals $\Psi$ is weakly lower semi-continuous.

Proof . For each $i=0, \ldots, N$ and any $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$, we can define the functionals $J_{i}$ and $H: W_{0}^{1, \vec{p}(\cdot)}(\Omega) \longrightarrow \mathbb{R}$ as follows:

$$
\begin{gathered}
J_{i}=\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x, \quad \text { where } \quad \frac{\partial u}{\partial x_{0}}=u \\
H(u)=-\int_{\Omega} G(x, u) d x
\end{gathered}
$$

Claim 1: Consider a sequence $\left(u_{n}\right)_{n}$ with the property that $u_{n} \rightharpoonup u$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. As $J_{i}$ is convex, for every $n$, we obtain

$$
J_{i}(u) \leq J_{i}\left(u_{n}\right)+\left\langle J_{i}^{\prime}(u), u-u_{n}\right\rangle .
$$

Taking the limit as $n \rightarrow \infty$ in the above inequality, we observe that $J_{i}$ is sequentially weakly lower semi-continuous. As a result, we obtain:

$$
\begin{equation*}
\int_{\Omega} \sum_{i=1}^{N} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x \leq \liminf _{n \rightarrow+\infty} \int_{\Omega} \sum_{i=0}^{N} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x \tag{3.9}
\end{equation*}
$$

By utilizing 3.9 and considering the continuity and monotonicity of $\widehat{M}_{i}$, we obtain

$$
\begin{align*}
\liminf _{n \rightarrow+\infty} J\left(u_{n}\right) & =\liminf _{n \rightarrow+\infty} \sum_{i=0}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq \sum_{i=0}^{N} \widehat{M}_{i}\left(\liminf _{n \rightarrow+\infty} \int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq \sum_{i=0}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq J(u) . \tag{3.10}
\end{align*}
$$

That is to say, $J$ demonstrates sequential weak lower semi-continuity.
Claim 2: $H$ is sequentially weakly continuous. Let $\left(u_{n}\right)_{n}$ be a sequence in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that $u_{n} \rightharpoonup u$ weakly in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. So, by (3.4) and Proposition 2.3. Therefore, it is easy to show that $\lim _{n \rightarrow \infty} H\left(u_{n}\right)=H(u)$, and hence $H$ is sequentially weakly lower semicontinuous. Similarly, just like we demonstrated for the mapping $H$, it is possible to establish the sequential weak lower semi-continuity of $\Phi$. Since $\Psi=J-H$, we complete the proof.

Lemma 3.5. On $K_{n}$, the functional $\Psi$ is boundedly below, and the infimum $m_{n}$ over $K_{n}$ is attained at $u_{n} \in K_{n}$.
Proof. To begin with, considering any $u \in K_{n}$, we find that

$$
\begin{align*}
\Psi(u) & =\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} G(x, u) d x \geq-\int_{\Omega} G(x, u) d x \\
& \geq-k b_{n} \operatorname{meas}(\Omega) \tag{3.11}
\end{align*}
$$

In conclusion, we deduce that $\Psi$ is bounded from below on $K_{n}$. It is apparent that $K_{n}$ possesses the properties of convexity and closedness, thus establishing its weak closedness within $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Consider the sequence $\left(u_{n}\right)_{n}$ in $K_{n}$ such that $\Psi\left(u_{n}\right)$ lies between $m_{n}$ and $m_{n}+\frac{1}{n}$ for all $n \in \mathbb{N}$, where $m_{n}=\inf _{K_{n}} \Psi$. Next, if $\left\|u_{n}\right\|_{1, \vec{p}(\cdot)} \leq 1$, our objective
is achieved; otherwise, we proceed with the following steps

$$
\begin{aligned}
\Psi(u) & =\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} G(x, u) d x \\
& =\sum_{i=1}^{N} \int_{0}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) M_{i}(s) d s-\int_{\Omega} G(x, u) d x \\
& \geq \frac{m}{\bar{p}} \sum_{i=1}^{N}\left(\left\|\frac{\partial u}{\partial x_{i}}\right\|_{p_{i}(\cdot)}^{\underline{p}}-1\right)-\int_{\Omega} G(x, u) d x \\
& \geq \frac{m}{\bar{p} N^{\underline{p}-1}}\left\|u_{n}\right\|_{1, \vec{p}(\cdot)}^{\underline{p}}-\frac{m N}{\bar{p} N^{\underline{p}-1}}-k b_{n}(\operatorname{meas} \Omega) .
\end{aligned}
$$

Which yields

$$
\begin{equation*}
\frac{m}{\bar{p} N^{\underline{p}-1}}\left\|u_{n}\right\|_{1, \vec{p}(\cdot)}^{\frac{p}{p}} \leq m_{n}+1+\frac{m N}{\bar{p} N^{\underline{p}-1}}+k b_{n}|\Omega| \tag{3.12}
\end{equation*}
$$

for all $n \in \mathbb{N}$, thus $\left(u_{n}\right)_{n}$ is bounded in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ which is a reflexive space. Therefore, by considering a sub-sequence denoted as $\left(u_{n}\right)_{n}$, we observe weak convergence towards a specific element $u_{n} \in K_{n}$. This leads us to the conclusion that $\Psi\left(u_{n}\right)=m_{n}$, utilizing the concept of weakly sequentially lower semi-continuity of $\Psi$.

Step 2 :. A priori estimates.
We start this step by proving in the following result that the sequence $\left(u_{n}\right)_{n}$ is bounded almost everywhere.
Proposition 3.6. For all $n \in \mathbb{N}$, we have $0 \leq u_{n}(x) \leq a_{n}$ a.e. $x \in \Omega$.
Proof . Let $\Lambda_{n}=\left\{x \in \Omega: b_{n} \geq u_{n}(x)>a_{n}\right\}$ and suppose that meas $\left(\Lambda_{n}\right)>0$. Define the function $\sigma_{n}(t)=$ $\min \left(\max (t, 0), a_{n}\right)$ and set $h_{n}=\sigma_{n}\left(u_{n}\right)$. It is clear that from the definition and the continuity of $\sigma_{n}$ we get $h_{n} \in K_{n}$. As a consequence, we obtain that

$$
h_{n}(x)= \begin{cases}u_{n}(x) & \text { if } x \in \Omega \backslash \Lambda_{n}  \tag{3.13}\\ a_{n} & \text { if } x \in \Lambda_{n}\end{cases}
$$

Then, we can write

$$
\begin{align*}
& \Psi\left(h_{n}\right)-\Psi\left(u_{n}\right) \\
= & \sum_{i=1}^{N} \widehat{M_{i}}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial h_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} \int_{0}^{h_{n}} g(x, t) d t d x-\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\int_{\Omega} \int_{0}^{u_{n}} g(x, t) d t d x \\
= & \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega \backslash \Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial h_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial h_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& -\int_{\Omega \backslash \Lambda_{n}} \int_{0}^{h_{n}} g(x, t) d t d x-\int_{\Lambda_{n}} \int_{0}^{h_{n}} g(x, t) d t d x \\
& -\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega \backslash \Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& +\int_{\Omega \backslash \Lambda_{n}} \int_{0}^{u_{n}} g(x, t) d t d x+\int_{\Lambda_{n}} \int_{0}^{u_{n}} g(x, t) d t d x \\
= & -\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Lambda_{n}} \int_{0}^{a_{n}} g(x, t) d t d x+\int_{\Lambda_{n}} \int_{0}^{u_{n}} g(x, t) d t d x \\
= & -\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Lambda_{n}} \int_{u_{n}}^{h_{n}} g(x, t) d t d x \\
= & -\sum_{i=1}^{N} \widehat{M_{i}}\left(\int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Lambda_{n}} \int_{u_{n}}^{a_{n}} g(x, t) d t d x \leq 0 . \tag{3.14}
\end{align*}
$$

Because $\int_{\Lambda_{n}} \int_{u_{n}}^{a_{n}} g(x, t) d t d x \geq 0$. Hence, $\Psi\left(h_{n}\right) \geq \Psi\left(u_{n}\right)=\inf _{K_{n}} \Psi$, then every term should be zero. In particular,

$$
\begin{equation*}
\sum_{i=1}^{N} \int_{\Lambda_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x=\int_{\Lambda_{n}}\left(G\left(x, a_{n}\right)-G\left(x, u_{n}\right)\right) d x \tag{3.15}
\end{equation*}
$$

Therefore, $\operatorname{meas}\left(\Lambda_{n}\right)=0$, which means $0 \leq u_{n}(x) \leq a_{n}$ almost every where $x \in \Omega$. $\qquad$ Next we show that the sequence $\left(u_{n}\right)_{n}$ formed of weak solutions of problem (3.5) as mentioned in the following result.

Proposition 3.7. The terms of $\left(u_{n}\right)_{n}$ are local minimum points of $\Psi$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.
Proof . Set $\Gamma_{n}=\left\{x \in \Omega: b_{n} \geq u(x)>a_{n}\right\}$. So, we have $\int_{\sigma_{n}(u)}^{u} g(x, t) d t=0$ for any $x \in \Omega \backslash \Gamma_{n}$. In the other hand, if $x \in \Gamma_{n}$, then one has the following three cases.
(a) If $u(x)<0$, then $\int_{\sigma_{n}(u)}^{u} g(x, t) d t=0$.
(b) If $a_{n}<u(x) \leq b_{n}$, then by (H2). $\int_{\sigma_{n}(u)}^{u} g(x, t) d t \leq 0$.
(c) If $b_{n}<u(x)$, then $\int_{\sigma_{n}(u)}^{u} g(x, t) d t=\int_{a_{n}}^{u} g(x, t) d t \leq \int_{a_{n}}^{u} k d t=k\left(u(x)-a_{n}\right)$, by (3.4).

Fix a real $\underline{p}_{, \infty}$ such that $\underline{p}_{, \infty}>q(x)+1>\bar{p}$ for every $x \in \Omega$, then the following constant is finite

$$
\lambda=\sup _{\mu \geq b_{n}} \frac{k\left(\mu-a_{n}\right)}{\left(\mu-a_{n}\right)^{q(x)+1}} .
$$

Then, for almost every where $x \in \Omega$, we have $\int_{\sigma_{n}(u)}^{u} g(x, t) d t \leq \lambda\left|\left(u(x)-\sigma_{n}(u(x))\right)\right|^{q(x)+1}$. Then, since when $\underline{p} \leq N$, the space is $W^{1, \vec{p}(\cdot)}(\Omega)$ compactly embedded in $L^{q(\cdot)+1}(\Omega)$ and continuously embedded in $C^{0}(\bar{\Omega})$ elsewhere, there is a positive constant $c$ such that

$$
\begin{equation*}
\int_{\Omega} \int_{\sigma_{n}(u)}^{u} g(x, t) d t d x \leq c^{q(x)+1} \lambda\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1} \tag{3.16}
\end{equation*}
$$

Therefore, we can write

$$
\begin{align*}
\Psi(u)-\Psi\left(\sigma_{n}(u)\right)= & \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} \int_{0}^{u} g(x, t) d t d x \\
& -\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial \sigma_{n}(u)}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\int_{\Omega} \int_{0}^{\sigma_{n}(u)} g(x, t) d t d x \\
= & \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Gamma_{n}} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Gamma_{n}} \int_{\sigma_{n}(u)}^{u} g(x, t) d t d x \\
= & \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}-\frac{\partial \sigma_{n}(u)}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} \int_{\sigma_{n}(u)}^{u} g(x, t) d t d x \\
\geq & \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}-\frac{\partial \sigma_{n}(u)}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\lambda c^{q(x)+1}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1} \quad(\text { by (3.16) }) \\
\geq & \frac{m}{p N^{p-1}}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{\bar{p}}-\lambda c^{q(x)+1}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1} \tag{3.17}
\end{align*}
$$

Since $\sigma_{n}(u) \in K_{n}$, we have $\Psi\left(\sigma_{n}(u)\right) \geq \Psi\left(u_{n}\right)$ and preserving the generality of our analysis, let's assume that $\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)} \leq 1$ cause we need small values of $\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}$. Then

$$
\begin{align*}
\Psi(u) & \geq \Psi\left(u_{n}\right)+\frac{m}{\bar{p}(N)^{\underline{p}-1}}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{\bar{p}}-\lambda c^{q(x)+1}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1} \\
& \geq \Psi\left(u_{n}\right)+\left(\frac{m}{\bar{p}(N)^{\bar{p}-1}}-\lambda c^{q(x)+1}\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1-\bar{p}}\right)\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{\bar{p}} . \tag{3.18}
\end{align*}
$$

The continuity of $\sigma_{n}$ allows us to choose a positive value $\delta>0$ such that, for any $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$, the condition

$$
\begin{equation*}
\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}<\delta, \quad\left\|\left(u-\sigma_{n}(u)\right)\right\|_{1, \vec{p}(\cdot)}^{q(x)+1-\bar{p}} \leq \frac{m}{\bar{p}(N)^{\underline{p}-1} \lambda c^{q(x)+1}}, \tag{3.19}
\end{equation*}
$$

this implies that $u_{n}$ is a local minimum of $\Psi$.
Proposition 3.8. The sequence $\left(m_{n}\right)_{n}$ is strictly negative and converges to zero.
Proof. In view of condition (H3), we have $\vartheta_{n} \in K_{n}$

$$
\begin{equation*}
m_{n} \leq \Psi\left(\vartheta_{n}\right)=-\int_{\Omega} \int_{0}^{\vartheta_{n}} f(x, t) d t d x<0 \tag{3.20}
\end{equation*}
$$

To prove that $\lim _{n \rightarrow+\infty} m_{n}=0$ it is sufficient to observe that for every $n \in \mathbb{N}$ and $u \in K_{n}$, we have

$$
\begin{equation*}
0>m_{n}=\Psi\left(u_{n}\right) \geq-k b_{n}|\Omega| . \tag{3.21}
\end{equation*}
$$

Since $\left(b_{n}\right)_{n}$ converges to zero, we conclude the required result.
Step 3 :. Proof of Theorem 3.2. Since the terms of $\left(u_{n}\right)_{n}$ are local minima of $\Psi$, they are weak solutions of (1.2). In virtue of Proposition 3.6 we have $0 \leq u_{n}(x) \leq a_{n}$ for almost every where $x \in \Omega$ and since $\left(a_{n}\right)_{n}$ converges to zero. An infinite number of distinct sequences $\left(u_{n}\right)_{n}$ can be found such that $\lim _{n \rightarrow+\infty}\left\|u_{n}\right\|_{L^{\infty}(\Omega)}=0$. Furthermore, we have

$$
\begin{align*}
m_{n}=\Psi\left(u_{n}\right) & \geq \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-\int_{\Omega} G\left(x, u_{n}\right) d x \\
& \geq \sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)-k b_{n}|\Omega| \tag{3.22}
\end{align*}
$$

Thus, if $\left\|u_{n}\right\|_{1, \vec{p}(\cdot)} \leq 1$, we have

$$
\begin{equation*}
\frac{m}{\bar{p} N^{\underline{p}-1}}\left\|u_{n}\right\|_{1, \vec{p}(\cdot)}^{\bar{p}} \leq m_{n}+k b_{n}|\Omega| \longrightarrow 0 . \tag{3.23}
\end{equation*}
$$

Thus, $\lim _{n \rightarrow+\infty}\|u\|_{1, \vec{p}(\cdot)}=0$, which completes our proof.
Now, we present an example to illustrate the main results.
Example 3.9. We define $M_{i}(t)=(1+t)^{\theta_{i}}$ for $i=1, \ldots N$, where $\theta_{i}>0$. It is worth noting that $M_{i}(t) \geq 1$ for all $t \geq 0$, which directly verifies the condition stated in (H4). Let

$$
f(x, t)=\left\{\begin{array}{lc}
\left(1+|x|^{2}\right)(\underline{p}+2) t^{\underline{p}}+1 \sin \left(\frac{1}{t^{\underline{p}}}\right)-\underline{p}\left(1+|x|^{2}\right) t \cos \left(\frac{1}{t^{\underline{p}}}\right) & \text { if } t>0  \tag{3.24}\\
0 & \text { otherwise }
\end{array}\right.
$$

It is easy to compute directly that

$$
F(x, t)=\left\{\begin{array}{lc}
\left(1+|x|^{2}\right) t^{\underline{p}+2} \sin \left(\frac{1}{t^{\underline{p}}}\right) & \text { if } t>0  \tag{3.25}\\
0 & \text { otherwise }
\end{array}\right.
$$

We shall now consider the following nonlinear perturbed Kirchhoff problem

$$
\left\{\begin{array}{l}
-\sum_{i=1}^{N}\left[1+\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right]^{\theta_{i}} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right)  \tag{3.26}\\
=\left(1+|x|^{2}\right)(\underline{p}+2) t^{\underline{p}}+1 \sin \left(\frac{1}{t^{\underline{\underline{p}}}}\right)-\underline{p}\left(1+|x|^{2}\right) t \cos \left(\frac{1}{t^{\underline{p}}}\right) \quad \text { in } \Omega, \\
u=0 \text { on } \bar{\partial} \Omega .
\end{array}\right.
$$

Let $\left(a_{n}\right)_{n},\left(b_{n}\right)_{n}$, and $\left(\vartheta_{n}\right)_{n}$ be three positive sequences satisfying the conditions:

$$
\begin{equation*}
a_{n}=\left(\frac{1}{2 n \pi+2 \pi}\right)^{\frac{1}{\underline{p}}}, \quad b_{n}=\left(\frac{1}{2 n \pi+\frac{3 \pi}{2}}\right)^{\frac{1}{\underline{p}}} \quad \text { and } \quad \vartheta_{n}=\left(\frac{1}{4 n \pi+\frac{\pi}{2}}\right)^{\frac{1}{\underline{p}}} \tag{3.27}
\end{equation*}
$$

for every $n \in \mathbb{N}$. Then one easily deduces

$$
\int_{0}^{a_{n}} f(x, s) d s=\sup _{t \in\left[a_{n}, b_{n}\right]} \int_{0}^{t} f(x, s) d s
$$

and $F\left(x, \vartheta_{n}\right)>0$. So conditions (H2) and (H3) have been verified. Having satisfied all the assumptions of Theorem 3.2 we can affirm the existence of a sequence of positive, homoclinic weak solutions $\left(u_{n}\right)_{n}$ in $W^{1, \vec{p}(x)}(\Omega)$ for the problem (3.26).

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