

# On the Fekete-Szegö problem associated with generalized fractional operator

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## Abstract

In this paper, the classical Fekete-Szegö problem is studied regarding a class of univalent functions generated using a generalized fractional differential operator. The results presented in the main theorem are new generalizations for well-known results.

**Keywords:** Analytic functions, fractional operator, univalent functions, normalized functions, Fekete-Szegö problems

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## 1 Introduction

We denote by  $A$  the functions in the unit disc  $U = \{z : 0 < |z| < 1\}$  that has the form

$$f(z) = z + \sum_{n=2}^{\infty} a_n z^n, \quad a_n \geq 0. \quad (1.1)$$

Let  $S$  denote the class of univalent functions in the unit disk. The Hadamard convolution of the functions  $f(z) = z + \sum_{n=2}^{\infty} a_n z^n$  and  $g(z) = z + \sum_{n=2}^{\infty} b_n z^n$  is given by:

$$f(z) * g(z) = z + \sum_{n=2}^{\infty} a_n b_n z^n.$$

For a class  $S$  of analytic functions, Fekete and Szegö [5] obtained sharp upper bound for  $|a_3 - \mu a_2^2|$ , when  $\mu$  is real. Later on, this inequality has received a huge interest from many researchers. For very recent results regarding Fekete-Szegö problem studied for different new classes of functions see the papers [2, 7, 8, 13, 16] and [17]. In this paper, we obtain sharp upper bound for the Fekete-Szegö inequality regarding a certain subclass of  $S$ , generated using the generalized fractional operator that has been introduced by Issa and Darus [9], given as follows:

$$D_z^{\nu, m} f(z) = z + \sum_{n=2}^{\infty} \frac{\Gamma(n+1) \Gamma(2-m\nu)}{\Gamma(n-m\nu+1)} a_n z^n,$$

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where, applications of this operator can be found on [8] and [10]. Now we define the class by  $\mathcal{M}^{\nu,m}(\Phi, \Psi; \lambda, \alpha, \beta)$ , as follows:

**Definition 1.1.** A function  $f(z) = z + \sum_{n=2}^{\infty} a_n z^n$ ,  $a_n \geq 0$  belongs to the class  $\mathcal{M}^{\nu,m}(\Phi, \Psi; \lambda, \alpha, \beta)$ ;  $0 \leq \alpha < 1, 0 \leq \lambda \leq 1, \beta > 0, 0 \leq \nu < 1$ , and  $m = 1, 2, \dots$ , if there exist a function  $g \in A$  such that:

$$\operatorname{Re} \left( \frac{\lambda z^2 (D_z^{\nu,m} f(z))'' + z (D_z^{\nu,m} f(z))'}{\lambda z g'(z) + (1 - \lambda) g(z)} \right) > \alpha, z \in U \quad (1.2)$$

where  $g(z) = z + b_2 z + b_3 z^2 + \dots$ , is analytic and satisfies

$$\left| \arg \left( \frac{g(z) * \Phi(z)}{g(z) * \Psi(z)} \right) \right| < \frac{\pi\beta}{2}, \quad (1.3)$$

for some  $\Phi(z) = z + \sum_{n=2}^{\infty} \varpi_n z^n$ , and  $\Psi(z) = z + \sum_{n=2}^{\infty} \gamma_n z^n$ , both are analytic in  $U$  such that  $g(z) * \Psi(z) \neq 0$ ,  $\varpi_n \geq 0, \gamma_n \geq 0$ , and  $\varpi_n \geq \gamma_n$  ( $n \geq 2$ ).

Note that  $\mathcal{M}^{\nu,0}(\Phi, \Psi; \lambda, \alpha, \beta)$  give us the class  $\mathcal{M}(\Phi, \Psi; \lambda, \alpha, \beta)$  that was defined by Darus in [4],

$$\mathcal{M}^{\nu,0} \left( \frac{z}{(1-z)^2}, \frac{z}{1-z}; 0, 0, \beta \right) = K(\beta),$$

the class of closed-to-convex that was defined by Chonweerayoot et al. in [1], and  $\mathcal{M}^{\nu,0} \left( \frac{z}{(1-z)^2}, \frac{z}{1-z}; 0, 0, 1 \right)$  is the class of normalized close-to-convex functions was defined by Kaplan in [12].

Furthermore,  $\mathcal{M}^{\nu,0} \left( \frac{z}{(1-z)^2}, \frac{z}{1-z}; 0, \alpha, \beta \right) = K(\alpha, \beta)$  is the class of normalized close-to-convex defined by Darus and Thomas in [3]. We have to mention that the result we introduce in the following section generalizes various results obtained by many researchers.

## 2 Main result

To get our main result we need the following definition and lemma:

**Definition 2.1.** The analytic function  $p(z)$  in  $U$  belongs to the class  $P$ , if  $p(z) = 1 + p_1 z + p_2 z^2 + p_3 z^3 + \dots$  and  $\operatorname{Re} p(z) > 0$  for all  $z$  in the open unit disk.

**Lemma 2.1.** [15] Let  $h(z) = 1 + c_1 z + c_2 z^2 + \dots$ , be analytic in the open unit disk with  $\operatorname{Re} h(z) > 0$  for all  $z \in U$ . Then

$$\left| c_2 - \frac{c_1^2}{2} \right| \leq 2 - \frac{|c_1|^2}{2}. \quad (2.1)$$

Now we present and prove our main result in the following theorem.

**Theorem 2.2.** Let  $f(z) = z + a_2 z^2 + a_3 z^3 + \dots$ , be a function in the class  $\mathcal{M}^{\nu,m}(\Phi, \Psi; \lambda, \alpha, \beta)$ ;  $0 \leq \alpha < 1, 0 \leq \lambda \leq 1, \beta > 0, 0 \leq \nu < 1$ , and  $m = 1, 2, \dots$ , If  $3\eta\mu \geq 2\delta^2 + 4\delta\gamma_2$ , where  $\delta = \varpi_2 - \gamma_2, \eta = \varpi_3 - \gamma_3$  and  $\mu \geq 1$ . Then we have the sharp inequality

$$\begin{aligned} |a_3 - \mu a_2^2| &\leq \frac{\beta^2}{3\delta^2\eta C_4(\nu, m) C_3^2(\nu, m)} \{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta + 2\gamma_2)\} \\ &\quad - \frac{3\delta\mu\alpha(1-\alpha)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m) C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\ &\quad + \frac{2\beta(1-\alpha)(1+\lambda)}{3\delta C_3^2(\nu, m) C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \{3\mu C_4(\nu, m) - 2C_3^2(\nu, m)(1+\lambda)^2\} \\ &\quad + \frac{(1-\alpha)}{3\delta C_3^2(\nu, m) C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \{3\mu\delta C_4(\nu, m)(1+2\lambda) - 2C_3^2(\nu, m)(1+\lambda)^2\}, \end{aligned}$$

where,  $C_3(\nu, m) = \frac{\Gamma(3)\Gamma(2-m\nu)}{\Gamma(3-m\nu)}$  and  $C_4(\nu, m) = \frac{\Gamma(4)\Gamma(2-m\nu)}{\Gamma(4-m\nu)}$ .

**Proof .** Since  $f(z) \in \mathcal{M}^{\nu, m}(\Phi, \Psi; \lambda, \alpha, \beta)$ , it follows from (1.2) that there exist  $q(z) = 1 + q_1 z + q_2 z^2 + \dots$ , with  $q(z) \in P$ , and  $g(z) = z + b_2 z^2 + b_3 z^3 + \dots \in A$ , such that

$$\lambda z^2 (D_z^{\nu, m} f(z))'' + z (D_z^{\nu, m} f(z))' = \{\alpha + (1 - \alpha) q(z)\} \left\{ \lambda z g'(z) + (1 - \lambda) g(z) \right\},$$

which implies

$$\begin{aligned} & z + 2C_3(\nu, m)(\lambda + 1)a_2 z^2 + 3C_4(\nu, m)(1 + 2\lambda)a_3 z^3 + \dots \\ &= z + \{(1 + \lambda)b_2 + (1 - \alpha)q_1\}z^2 + \{(1 + 2\lambda)b_3 + (1 - \alpha)(1 + \lambda)q_1 b_2 + (1 - \alpha)q_2\}z^3 + \dots \end{aligned}$$

Equating the coefficients in the last equation implies

$$2C_3(\nu, m)(1 + \lambda)a_2 = (1 + \lambda)b_2 + (1 - \alpha)q_1, \quad (2.2)$$

and

$$3C_4(\nu, m)(1 + 2\lambda)a_3 = (1 + 2\lambda)b_3 + (1 - \alpha)(1 + \lambda)q_1 b_2 + (1 - \alpha)q_2. \quad (2.3)$$

Moreover, from (1.3), there exist a function in  $P$  say,  $p(z) = 1 + p_1 z + p_2 z^2 + p_3 z^3 + \dots$ , such that

$$g(z) * \Phi(z) = (g(z) * \Psi(z)) p^\beta(z),$$

which implies

$$\begin{aligned} & (z + b_2 z^2 + b_3 z^3 + \dots) * (z + \varpi_2 z^2 + \varpi_3 z^3 + \dots) \\ &= \{(z + b_2 z^2 + b_3 z^3 + \dots) * (z + \gamma_2 z^2 + \gamma_3 z^3 + \dots)\} (1 + p_1 z + p_2 z^2 + p_3 z^3 + \dots)^\beta. \end{aligned}$$

So,

$$\begin{aligned} & z + b_2 \varpi_2 z^2 + b_3 \varpi_3 z^3 + \dots = \{z + b_2 \gamma_2 z^2 + b_3 \gamma_3 z^3 + \dots\} \\ & \quad \left\{ 1 + \beta p_1 z + \left( \beta p_2 + \left( \frac{\beta^2}{2} - \frac{\beta}{2} \right) p_1^2 \right) z^2 + \dots \right\} \\ &= z + (p_1 \beta + \gamma_2 b_2) z^2 \\ & \quad + \left( p_1 \beta b_2 \gamma_2 + \beta p_2 + \left( \frac{\beta^2}{2} - \frac{\beta}{2} \right) p_1^2 + \gamma_3 b_3 \right) z^3 + \dots \end{aligned}$$

Equating the coefficients in the last equality implies

$$b_2 \varpi_2 = p_1 \beta + \gamma_2 b_2,$$

which give us

$$b_2 = \frac{p_1 \beta}{\delta}. \quad (2.4)$$

Also,

$$\begin{aligned} b_3 \varpi_3 &= p_1 \beta b_2 \gamma_2 + \beta p_2 + \left( \frac{\beta^2}{2} - \frac{\beta}{2} \right) p_1^2 + \gamma_3 b_3 \\ b_3 \varpi_3 - \gamma_3 b_3 &= p_1 \beta b_2 \gamma_2 + \beta p_2 + \frac{\beta^2}{2} p_1^2 - \frac{\beta}{2} p_1^2 \\ (\varpi_3 - \gamma_3) b_3 &= \beta \left( p_1 b_2 \gamma_2 + p_2 + \frac{\beta}{2} p_1^2 - \frac{1}{2} p_1^2 \right) \end{aligned}$$

$$\begin{aligned}\eta b_3 &= \beta \left( p_1 \frac{p_1 \beta}{\delta} \gamma_2 + p_2 + \frac{\beta}{2} p_1^2 - \frac{1}{2} p_1^2 \right) = \beta \left( p_2 + p_1^2 \frac{\beta}{\delta} \gamma_2 + \frac{\beta}{2} p_1^2 - \frac{1}{2} p_1^2 \right) \\ &= \beta \left( p_2 + p_1^2 \left( \frac{\beta}{\delta} \gamma_2 + \frac{\beta}{2} - \frac{1}{2} \right) \right) = \beta \left( p_2 + p_1^2 \left( \frac{2\beta\gamma_2}{2\delta} + \frac{\delta\beta}{2\delta} - \frac{\delta}{2\delta} \right) \right),\end{aligned}$$

which implies

$$b_3 = \frac{\beta}{\eta} \left\{ p_2 + \left( \frac{\beta(\delta + 2\gamma_2) - \delta}{2\delta} \right) p_1^2 \right\}. \quad (2.5)$$

So from (2.1),(2.2),(2.3),(2.4) and (2.5) we get

$$a_2 = \frac{p_1 \beta}{2\delta C_3(\nu, m)} + \frac{(1-\alpha) q_1}{2(1+\lambda) C_3(\nu, m)}, \quad (2.6)$$

and

$$\begin{aligned}a_3 &= \frac{1}{3C_4(\nu, m)(1+2\lambda)} \{ (1+2\lambda) b_3 + (1-\alpha)(1+\lambda) q_1 b_2 + (1-\alpha) q_2 \} \\ &= \frac{(1+2\lambda) b_3}{3C_4(\nu, m)(1+2\lambda)} + \frac{(1-\alpha)(1+\lambda) q_1 b_2}{3C_4(\nu, m)(1+2\lambda)} + \frac{(1-\alpha) q_2}{3C_4(\nu, m)(1+2\lambda)}.\end{aligned}$$

Indeed,

$$a_3 = \frac{\beta}{3\eta C_4(\nu, m)} \left\{ p_2 + \left( \frac{\beta(\delta + 2\gamma_2) - \delta}{2\delta} \right) p_1^2 \right\} + \frac{(1-\alpha)(1+\lambda) q_1 p_1 \beta}{3\delta C_4(\nu, m)(1+2\lambda)} + \frac{(1-\alpha) q_2}{3C_4(\nu, m)(1+2\lambda)}. \quad (2.7)$$

Now, we will use (2.1),(2.2),(2.3),(2.4),(2.5),(2.6) and (2.7) to get the required result as follows:

$$\begin{aligned}a_3 - \mu a_2^2 &= \frac{\beta p_2}{3\eta C_4(\nu, m)} + \frac{\beta^2 (\delta + 2\gamma_2) p_1^2}{6\delta\eta C_4(\nu, m)} - \frac{\beta \delta p_1^2}{6\delta\eta C_4(\nu, m)} + \frac{(1-\alpha)(1+\lambda) q_1 p_1 \beta}{3\delta C_4(\nu, m)(1+2\lambda)} + \frac{(1-\alpha) q_2}{3C_4(\nu, m)(1+2\lambda)} \\ &\quad - \frac{\mu \beta^2 p_1^2}{4\delta^2 C_3^2(\nu, m)} - \frac{2\beta(1-\alpha) \mu q_1 p_1}{4\delta C_3^2(\nu, m)(1+\lambda)} - \frac{(1-\alpha)^2 \mu q_1^2}{4(1+\lambda)^2 C_3^2(\nu, m)} \\ &\quad + \frac{2(1-\alpha) q_1^2}{12C_4(\nu, m)(1+2\lambda)} - \frac{2(1-\alpha) q_1^2}{12C_4(\nu, m)(1+2\lambda)} \\ &= \left\{ \frac{(1-\alpha) q_2}{3C_4(\nu, m)(1+2\lambda)} - \frac{2(1-\alpha) q_1^2}{12C_4(\nu, m)(1+2\lambda)} \right\} + \left\{ \frac{\beta p_2}{3\eta C_4(\nu, m)} - \frac{\beta \delta p_1^2}{6\delta\eta C_4(\nu, m)} \right\} \\ &\quad + \left\{ \frac{2(1-\alpha) q_1^2}{12C_4(\nu, m)(1+2\lambda)} - \frac{(1-\alpha)^2 \mu q_1^2}{4(1+\lambda)^2 C_3^2(\nu, m)} \right\} + \left\{ \frac{(1-\alpha)(1+\lambda) q_1 p_1 \beta}{3\delta C_4(\nu, m)(1+2\lambda)} - \frac{2\beta(1-\alpha) \mu q_1 p_1}{4\delta C_3^2(\nu, m)(1+\lambda)} \right\} \\ &\quad + \left\{ \frac{\beta^2 (\delta + 2\gamma_2) p_1^2}{6\delta\eta C_4(\nu, m)} - \frac{\mu \beta^2 p_1^2}{4\delta^2 C_3^2(\nu, m)} \right\},\end{aligned}$$

implies

$$\begin{aligned}a_3 - \mu a_2^2 &= \frac{(1-\alpha)}{3C_4(\nu, m)(1+2\lambda)} \left\{ q_2 - \frac{q_1^2}{2} \right\} + \frac{\beta}{3\eta C_4(\nu, m)} \left\{ p_2 - \frac{p_1^2}{2} \right\} \\ &\quad + (1-\alpha) q_1^2 \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3\mu(1-\alpha) C_4(\nu, m)(1+2\lambda)}{12C_3^2(\nu, m) C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} \\ &\quad + (1-\alpha) q_1 p_1 \beta \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3C_4(\nu, m) \mu(1+2\lambda)}{6\delta C_3^2(\nu, m) C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\} \\ &\quad + \beta^2 p_1^2 \left\{ \frac{2C_3^2(\nu, m) \delta(\delta + 2\gamma_2) - 3\mu\eta C_4(\nu, m)}{12\delta^2\eta C_4(\nu, m) C_3^2(\nu, m)} \right\}.\end{aligned} \quad (2.8)$$

Secondly, we will complete the proof by applying the technique introduced by London [14]. Now, assume that  $a_3 - \mu a_2^2$  is positive, then from (2.8) and by substituting  $p_1 = 2re^{i\theta}$  and  $q_1 = 2Re^{i\varphi}, 0 \leq r \leq 1, 0 \leq R \leq 1, 0 \leq \theta \leq 2\pi$ , and  $0 \leq \varphi \leq 2\pi$ , we obtain:

$$\begin{aligned} 3 \operatorname{Re}(a_3 - \mu a_2^2) &= \frac{(1-\alpha)}{C_4(\nu, m)(1+2\lambda)} \operatorname{Re} \left\{ q_2 - \frac{q_1^2}{2} \right\} + \frac{\beta}{\eta C_4(\nu, m)} \operatorname{Re} \left\{ p_2 - \frac{p_1^2}{2} \right\} \\ &\quad + (1-\alpha) \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)}{4C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} \operatorname{Re} q_1^2 \\ &\quad + \beta(1-\alpha) \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3C_4(\nu, m)\mu(1+2\lambda)}{2\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\} \operatorname{Re} p_1 q_1 \\ &\quad + \beta^2 \left\{ \frac{2C_3^2(\nu, m)\delta(\delta+2\gamma_2) - 3\mu\eta C_4(\nu, m)}{4\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} \operatorname{Re} p_1^2. \end{aligned}$$

Consequently, by Lemma 2.1 we get

$$\begin{aligned} 3 \operatorname{Re}(a_3 - \mu a_2^2) &\leq \frac{2(1-\alpha)}{C_4(\nu, m)(1+2\lambda)} (1-R^2) + \frac{2\beta}{\eta C_4(\nu, m)} (1-r^2) \\ &\quad + (1-\alpha) \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)}{4C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} 4R^2 \cos 2\theta \\ &\quad + \beta(1-\alpha) \left\{ \frac{2(1+\lambda)^2 C_3^2(\nu, m) - 3C_4(\nu, m)\mu(1+2\lambda)}{2\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\} 4rR \cos(\varphi + \theta) \\ &\quad + \beta^2 \left\{ \frac{2C_3^2(\nu, m)\delta(\delta+2\gamma_2) - 3\mu\eta C_4(\nu, m)}{4\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} 4r^2 \cos 2\varphi \\ &\leq \frac{2(1-\alpha)}{C_4(\nu, m)(1+2\lambda)} (1-R^2) + \frac{2\beta}{\eta C_4(\nu, m)} (1-r^2) \\ &\quad + (1-\alpha) \left\{ \frac{3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda) - 2(1+\lambda)^2 C_3^2(\nu, m)}{C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} R^2 \\ &\quad + 2\beta(1-\alpha) \left\{ \frac{3C_4(\nu, m)\mu(1+2\lambda) - 2(1+\lambda)^2 C_3^2(\nu, m)}{\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\} rR \\ &\quad + \beta^2 \left\{ \frac{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)}{\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} r^2 \\ &= H(r, R). \end{aligned}$$

Next, we will identify the sign of  $H_{rr}H_{RR} - (H_{rR})^2$  to illustrate the maximum value of  $H(r, R)$ , as follows:

$$\begin{aligned} H_{rr}H_{RR} - (H_{rR})^2 &= \left\{ \frac{-4\beta}{\eta C_4(\nu, m)} + 2\beta^2 \left\{ \frac{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)}{\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} \right\} \\ &\quad \left\{ \begin{aligned} &\frac{-4(1-\alpha)}{C_4(\nu, m)(1+2\lambda)} \\ &+ 2(1-\alpha) \left\{ \frac{3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda) - 2(1+\lambda)^2 C_3^2(\nu, m)}{C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} \\ &- \left\{ 2\beta(1-\alpha) \frac{3C_4(\nu, m)\mu(1+2\lambda) - 2(1+\lambda)^2 C_3^2(\nu, m)}{\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\}^2 \end{aligned} \right\} \\ &= \frac{16\beta(1-\alpha)}{\eta(1+2\lambda)C_4^2(\nu, m)} - 8\beta(1-\alpha) \left\{ \frac{3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda) - 2(1+\lambda)^2 C_3^2(\nu, m)}{C_3^2(\nu, m)C_4^2(\nu, m)\eta(1+2\lambda)(1+\lambda)^2} \right\} \end{aligned}$$

$$\begin{aligned}
& -8(1-\alpha)\beta^2 \left\{ \frac{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)}{\delta^2\eta C_4^2(\nu, m)C_3^2(\nu, m)(1+2\lambda)} \right\} \\
& + 4\beta^2(1-\alpha) \left[ \left\{ \frac{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)}{\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} \right. \\
& \quad \left. \left\{ \frac{3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda) - 2(1+\lambda)^2C_3^2(\nu, m)}{C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} \right] \\
& - 4\beta^2(1-\alpha)^2 \left\{ \frac{3C_4(\nu, m)\mu(1+2\lambda) - 2(1+\lambda)^2C_3^2(\nu, m)}{\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\}^2 \\
& < 0.
\end{aligned}$$

Hence the function  $H(r, R)$  attains its maximum value on the boundary of the unit disk. Therefore our inequality follows by observing that

$$\begin{aligned}
|a_3 - \mu a_2^2| & \leq \frac{1}{3}H(r, R) \leq \frac{1}{3}H(1, 1) \leq \beta^2 \left\{ \frac{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \right\} \\
& + (1-\alpha) \left\{ \frac{3\mu(1-\alpha)C_4(\nu, m)(1+2\lambda) - 2(1+\lambda)^2C_3^2(\nu, m)}{3C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \right\} \\
& + 2\beta(1-\alpha) \left\{ \frac{3C_4(\nu, m)\mu(1+2\lambda) - 2(1+\lambda)^2C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)} \right\} \\
& = \frac{\beta^2}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)\} \\
& + \frac{6\beta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)(1+\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{4\beta(1-\alpha)(1+\lambda)^3C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& + \frac{3\delta\mu(1-\alpha)^2C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{2\delta(1-\alpha)(1+\lambda)^2C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& = \frac{\beta^2}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)\} \\
& + \frac{6\beta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)(1+\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{4\beta(1-\alpha)(1+\lambda)^3C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& + \frac{3\delta\mu(1-2\alpha+\alpha^2)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{2\delta(1-\alpha)(1+\lambda)^2C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& = \frac{\beta^2}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)\} \\
& + \frac{6\beta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)(1+\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{4\beta(1-\alpha)(1+\lambda)^3C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& + \frac{3\delta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} + \frac{3\delta\mu\alpha(\alpha-1)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& - \frac{2\delta(1-\alpha)(1+\lambda)^2C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& = \frac{\beta^2}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \{3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta+2\gamma_2)\} - \frac{3\delta\mu\alpha(1-\alpha)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& + \frac{6\beta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)(1+\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{4\beta(1-\alpha)(1+\lambda)^3C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
& + \frac{3\delta\mu(1-\alpha)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} - \frac{2\delta(1-\alpha)(1+\lambda)^2C_3^2(\nu, m)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2}.
\end{aligned}$$

Finally, we get:

$$\begin{aligned}
 |a_3 - \mu a_2^2| &\leq \frac{\beta^2}{3\delta^2\eta C_4(\nu, m)C_3^2(\nu, m)} \left\{ 3\mu\eta C_4(\nu, m) - 2C_3^2(\nu, m)\delta(\delta + 2\gamma_2) \right\} \\
 &\quad - \frac{3\delta\mu\alpha(1-\alpha)C_4(\nu, m)(1+2\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \\
 &\quad + \frac{2\beta(1-\alpha)(1+\lambda)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \left\{ 3\mu C_4(\nu, m) - 2C_3^2(\nu, m)(1+\lambda)^2 \right\} \\
 &\quad + \frac{(1-\alpha)}{3\delta C_3^2(\nu, m)C_4(\nu, m)(1+2\lambda)(1+\lambda)^2} \left\{ 3\mu\delta C_4(\nu, m)(1+2\lambda) - 2C_3^2(\nu, m)(1+\lambda)^2 \right\},
 \end{aligned}$$

as required. The equality is attained by choosing  $\nu = 0$ ,  $p_1 = q_1 = 2i$  and  $p_2 = q_2 = -2$  in (2.8).  $\square$

**Remark 2.1.** Letting  $\nu = 0$  in Theorem 2.2, we have the result by Darus [4].

**Remark 2.2.** Letting  $\nu = 0$  and  $\lambda = 0$  in Theorem 2.2, we have the result by Frasin and Darus [6].

**Remark 2.3.** Letting  $\Phi(z) = \frac{z}{(1-z)^2}$ ,  $\Psi(z) = \frac{z}{1-z}$ ,  $\lambda = \alpha = 1$ , and  $\nu = 0$  in Theorem 2.2, we have the result by Jahangiri [11].

Finally, we ought to mention that we can gain a lot of results from Theorem 2.2 by various choices of  $\Phi(z)$  and  $\Psi(z)$ .

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